Before diving into rep'n theory of f.d. s.s. Lie algebras, let's necall some facts about the Lie algebras themselves. let L be a f.d. s.s. Lie algebra lover C or some alg. closed field k). Then L has a maximal total subalg. Fix one such algebra  $T \subset L$ , then  $L = T \oplus \bigoplus L_X$ where  $I = \{ X \in T^* \setminus \{ 0 \} : L_X \neq 0 \}$  is the set whose elements are called roots, and  $L_X = \{ n \in L : [t, n] = X(t) \times \forall t \in T \}$ Note.  $[L_X, L_p] \leq L_X + p$ .

Some facts about roots.  
(1) let 
$$\langle -, - \rangle$$
 denote the (dual) killing form on T\*.  
 $\forall v \in T^* \setminus \{0\}$ , let  $v' = \frac{2v}{\langle v, v \rangle}$ , then  
 $\langle \mu, v' \rangle = \frac{2 \langle \mu, v \rangle}{\langle v, v \rangle} = \frac{2}{\|v\|^2} \langle \mu, v \rangle$ 

The new bracket  $\langle -, -^{\vee} \rangle$  is linear only in the 1st variable and is intensitive to rescaling the inner polt  $\langle -, -\rangle$ . With respect to this new bracket,  $\forall d, p \in \overline{P}, \quad \langle d, p^{\vee} \rangle \in \overline{Z}$  $\subseteq Cartan integers$  Furthermore, (.)  $kd \cap \overline{\Psi} = \{\pm d\}$ (.)  $\forall a, \beta \in \overline{\Phi}, a - \langle a, \beta' \rangle \beta \in \overline{\Phi}$ (2)  $(sl_2 - triples)$ . Let  $x \in \overline{\Sigma}$ . Then  $\dim_{\mathbb{R}} l_{x} = 1$  and [Ld, L-2] = khd, where  $h_{d} \in T$  s.t.  $\langle \beta, d' \rangle = \beta(h_{d})$  $(Explicitly, h_d = \frac{2 t_d}{d(t_d)}, where t_d \in T \iff d \in T^*)$ Moreover,

 $s_{\lambda} = L_{-\lambda} \oplus kh \oplus L_{\lambda}$  is a lie subalg of  $L s.t. s_{\lambda} \cong sl_{2}$ , with h, a hesly.

(3) There is a finite set  $\Delta = \{ d_1, ..., d_n \}$  of roots that forms a basis for T\* = k" and s.t. YBE I,

B = E tod w/ all to E Z+ or all to E - Z+. The roots in  $\Delta$  are called simple roots. They give rise to a partition of I as > regative roots  $\overline{\Psi} = \overline{\Psi}_{+} \cup \overline{\Psi}_{-}$ positive mots

There is a partial order on T\* given by  $\mu \leq \nu \Leftrightarrow \nu - \mu \in \mathbb{Z}_{+} = \mathbb{Z}_{+} \Delta = \oplus \mathbb{Z}_{+} \alpha_{i}$ With this partial order, we can unite  $\overline{\Psi}_{+} = \{ \mathcal{A} \in \overline{\Psi} \mid \mathcal{A} > 0 \}$ 重-= { x e 重 | x × 0 } The root space decomposition takes the form  $N \pm = \bigoplus_{x \in \Phi+} L x$  $L = N_{-} \oplus T \oplus N_{+}$ 

where T is abelian and  $N_{\pm}$  are nilpotent. This is called the mangular decomposition of L. For the classical lie algs,  $N_{\pm}$  consist of shirtly upper/lower triangular matrices. We will also define the positive and negative Borel rubalg to be

$$B \pm = T \oplus N \pm .$$

(4) Fix a basis  $\Delta = \{ \alpha_1, ..., \alpha_n \}$  of simple roots of  $\mathbf{I}$ . Define the root lattice

$$R = \mathbb{Z} \overline{P} = \sum_{x \in \overline{P}} \mathbb{Z} d = \bigoplus_{i=1}^{n} \mathbb{Z} d_i \cong \mathbb{Z}^{\oplus n}$$

and the weight lattice

$$\Delta = \{ \lambda \in T^* | \langle \lambda, \alpha' \rangle \in \mathbb{Z} \; \forall \alpha \in \mathbb{P} \}$$

$$= \{ \lambda \in T^* | \langle \lambda, \alpha' \rangle \in \mathbb{Z} \; \forall \alpha \in \Delta \}$$

$$\int_{\Omega} vot obvious ! Exercise 7.2.3 in Lorenz's book$$

We have  $R \leq 1$ 

The weights  $\lambda_i$  s.t.  $\langle \lambda_i, d_j \rangle = \delta_{i,j}$  are called the fundamental weights. They form a Z-basis for  $\Lambda$ , i.e.  $\Lambda = \bigoplus_{i=1}^{\infty} Z \lambda_i \cong Z^{\otimes n}$ 

Next, we define

$$\Delta t = \left\{ \begin{array}{l} \lambda \in T^{*} \\ = \end{array} \right\} \langle \lambda, d_{i}^{*} \rangle \in \mathbb{Z} + \forall d_{i} \in \Delta \left\}$$

$$= \bigoplus_{i=1}^{n} \mathbb{Z} + \lambda_{i}$$

The weights  $\lambda \in A_{+}$  are called the dominant weights

THE WEYL GROUP

Let 
$$\overline{\Phi}$$
 be the set of roots and fix a basis  $\Delta$ .  
For each  $\chi \in \overline{\Phi}$ , let  $s_{\chi} \colon T^{\star} \to T^{\star}$   
 $S_{\chi}(\mu) = \mu - \langle \mu, \chi^{\vee} \rangle \chi$   
Can check that  $s_{\chi}$  is the inflection through the hyperplane  
orthogonal to  $\chi$ .

Let 
$$W = \langle S_{\alpha} | \alpha \in \Phi \rangle \subseteq GL(T^*)$$
  
 $W$  is called the Weyl group of  $\Phi$ .

Example (sl2)  
Recall that 
$$sl_2 = kf \oplus kh \oplus ke$$
  
 $f = \begin{pmatrix} 0 & 0 \\ 0 & -1 \end{pmatrix}, \quad e = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$   
 $[h, e] = de \quad [h, f] = -2f \qquad [e, f] = h$   
 $kh = a \text{ maximal total subalg}$   
 $(kh)^* \cong k, ie. \text{ functionals on } kh \text{ act as scalars}$   
With this identification,  $\overline{P} = \{ \pm 2 \}$  since  
 $ke = l_2 = f \times e sl_2 | Lh, \pi] = d\pi \{ \}$   
 $kf = l_{-2} = \{ \times e sl_2 | Lh, \pi] = -d\pi \}$ 



The not lattice is 22. Choose 
$$\Delta = \{2\}$$
, so  $\overline{\Psi}_{+} = \{2\}$   
and  $\overline{\Psi}_{-} = \{-2\}$ . Hence  $N_{+} = L_{2} = ke$   
 $N_{-} = L_{-2} = kf$   
The fund weight  $\lambda$  satisfies  $\langle \lambda, 2^{\vee} \rangle = 1$   
 $\langle \lambda, \frac{2}{4} 2 \rangle = \langle \lambda, 1 \rangle$   
so  $\lambda = 1$ . The weight lattice is 2. The weyle group  $\mathcal{W}$   
is generated by nellection  $s_{2}$  of order 2 so  $\mathcal{W} \cong 2/22$ ,  
and operates on  $T^{*} = k$  by multiplication by  $\pm 1$ .  
The partial order  $\leq m$   $T^{*} = k$  is given by  
 $\mu \leq \lambda \iff \lambda - \mu \in 22t$ .

let L be a f.d. s.s. Lie algebra. If 
$$V \in \operatorname{Rep}(L)^{fd}$$
,  
then we have a decomposition  
 $V \cong \bigoplus V_{\lambda}$  (X)  
 $\chi \in T^*$   
where  $V_{\lambda} = \{ \chi \in V : t. \chi = \lambda(t) \mid \chi \forall t \in T \}$  is the  $\lambda$ -  
weight space for V.  
Easy check :  $L_{X} \cdot V_{\lambda} \subseteq V_{X+\lambda}$ .

General result. (a) The weights 
$$\lambda$$
 occuring in (\*) belong  
to  $\Delta$ , is they appear in the weight lattice.  
(b) If  $\lambda$  is a weight of V, then the orbit  $W\lambda$  consist  
of weights of V, all having the same multiplicity:  
dimp  $V_{\lambda} = \dim_k V_{W\lambda}$   $\forall w \in W$ .

We will use this to calculate all weights appearing in certain f.d. representations later.

$$\frac{E \times ample}{V} (Rep(sl_2))$$
Let V be a follower of sl\_2 of dimil n, and consider
$$V \equiv \bigoplus_{\lambda \in T^*} V_{\lambda} \cong \bigoplus_{\lambda \in k} V_{\lambda}$$

Note (ke) 
$$\cdot V_{\lambda} = L_2 \cdot V_{\lambda} \subseteq V_{\lambda+2}$$
  
(kf)  $\cdot V_{\lambda} = L_{-2} \cdot V_{\lambda} \subseteq V_{\lambda-2}$ 

In particular, if  $\lambda$  is maximal among all weights wrt  $\leq$ , and  $0 \neq v \in V_{\lambda}$ , then  $e \cdot v = 0$  and v generates all of V. thence the vectors  $f^i \cdot v$ ,  $0 \leq i \leq n-1$ , forms a basis for V. Using this, we can figure out that the weights are n-1, n-3, ..., -(n-3), -(n-1)all occurring with multiplicity 1. A class of representatives of f.d. imaps are given by  $\operatorname{Sym}^m(k^2)$  where  $k^2$  is the natural rep'n,  $\forall m \geq 1$ 

Note that in this case, the finite dirich irreps only depend on the maximal weight. This remains the for general f.d.s.s. Lie algebras.

let  $V \in \text{Rep}(L)^{\text{fd}}$ . We say that V is a highest weight representation (with highest weight  $\lambda$ ) if X is a maximal weight (wrt  $\leq$ ) of all weights occuring in the weight decomposition  $V = \bigoplus V_{\mu}$ , and  $\exists 0 \neq v \in V_{\chi}$  e.t. vgenerates all of V (such a v is called a maximal vector).

A class of representatives of f.d. irreps is given by the Verma modules  $V(\lambda)$ ,  $\lambda \in \Lambda_+$ .

Some info about V(X):

(1) To compute which weight appears in 
$$V(\lambda)$$
, first  
compute all dominant weights  $\leq \lambda$ , then use Weyl group  
to get the rest. (lorenz' Proposition 7.16).  
(a) The multiplicity  $m_{\lambda}(\mu)$  of a weight  $\mu$  is given  
by Koctant's multiplicity formula  
 $m_{\lambda}(\mu) = \sum_{e \in V} (-1)^{e(e)} \cup (e(\lambda t p) - (\mu t p))$   
Kostant's partition function  
(3) The dimension of  $V(\lambda)$  is given by the Weyl dim

fremula: dim  $(V(\lambda)) = \prod_{\alpha \neq 0} \frac{\langle \lambda \neq \rho, \alpha \rangle}{\langle \rho, \alpha \rangle}$ .

(4) The multiplicity 
$$M_{\lambda, \mathcal{T}}^{M}$$
 of  $V(\mu)$  in  $V(\lambda) \otimes V(\mathcal{T})$   
is given by the Racah-Speicer formula  
 $M_{\lambda, \mathcal{T}}^{M} = \sum_{\sigma \in \mathcal{W}} (-1)^{e(\sigma)} m_{\mathcal{T}} (\sigma(\mu t \rho) - (\lambda t \rho)).$ 

Example From the Dynkin diagram of B2  
we have 2 simple roots 
$$d_1$$
,  $d_2$  w/ ||  $d_1$ || > ||  $d_2$  ||  
One computes angle  $(d_1, d_2) = 135$  degrees  
and  $\frac{||d_1||}{||d_2||} = \sqrt{2}$ .  
and can verify that the fund weights are  
 $\lambda_1 = \frac{\alpha_1}{2} + d_2$ ,  $\lambda_2 = d_1 + d_2$ .



Since  $2\lambda_2 \in \Lambda_1$ , we can consider  $V = V(2\lambda_2)$ . The dominant weights  $\leq 2\lambda_2$  are  $2\lambda_2$ ,  $\lambda_1$ ,  $\lambda_2$ , 0, which are rectangular nodes whose multiplicities inside are computed using (2). The remainder of the weight spaces in circular nodes are computed using the Weyl group.